Mini-Symposium Wuppertal

Schur-Weyl duality in molecular physics: A new glance at nuclear spin states in molecules

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December 2016

Conventional approximations in molecular theory

1st step Born-Oppenheimer approximation

 \Rightarrow Separate nuclear and electronic degrees of freedom

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> Molecular Hamiltonian and wave function $\mathcal{H} = \mathcal{H}_{el} + \mathcal{H}_{rovib} + \mathcal{H}_{nspin}$ $\Psi = \psi_{el} \ \psi_{rovib} \ \psi_{nspin}$

The exchange principle and permutation groups

Exchange of identical fermions \Rightarrow Sign change in wave function Exchange of identical **bosons** \Rightarrow No sign change

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The nuclear wave function

Molecular Hamiltonian and wave function

 $\mathcal{H} = \mathcal{H}_{\rm el} + \mathcal{H}_{\rm rovib} + \mathcal{H}_{\textit{nspin}}$

 $\Psi = \psi_{\rm el} \; \psi_{\rm rovib} \; \psi_{\rm nspin}$

The total nuclear wave function $\psi_{\text{rovib}} \psi_{\text{nspin}}$ must be (anti-)symmetric for (fermionic) bosonic nuclei.

Ro-vibrational part: Quantum numbers (J, ν_i) Nuclear spin part: # of identical particles and their spin

$$\psi_{\text{rovib}} \psi_{\text{nspin}} = \psi_{J,\nu_i} \times \begin{array}{c} \psi_{\text{nspin}}^1 \\ \vdots \\ \psi_{\text{nspin}}^g \end{array} \right\} \Rightarrow \Gamma_{\text{(anti)sym}}$$

g is the nuclear spin statistical weight

How to find the permutation symmetry of nuclear spin states

The pedestrian way

- Construct all possible combinations of individual spin states
- Determine their total spin and permutation symmetry

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Configuration	S ₂	$\mathrm{I}_{\mathrm{tot}}$	$M_{ m I}$
$\uparrow\uparrow$	A	1	1
$\downarrow\downarrow$	A	1	-1
$\uparrow\downarrow + \downarrow\uparrow$	A	1	0
$\uparrow\downarrow-\downarrow\uparrow$	В	0	0

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Example: The L' lengthy to						
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If yes, can I use it to simplify the calculation of the symmetry properties?

Preparation of the answer Unitary symmetry of spins

Assumption No influence of nuclear spin on molecular energy levels Implication Symmetry group of ψ_{nspin} : Unitary group U(d), d = 2I + 1

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Unitary group: The most basic transformations in quantum mechanics

- Restriction for symmetry of nuclear spin states: $|\langle \psi_{nspin} | \psi_{nspin} \rangle|^2$ conserved!
- All unitary transformations $U^{\dagger}U = \mathbb{1}$ fulfill this constraint.
- Dimension: 2I + 1 functions for single I

Unitary symmetry group: The irreducible representations

- U(d) Partition $\lambda = \{\lambda_1, \lambda_2, \dots, \lambda_d\}$ of positive descending integers is unambigous characterization of U(d) irrep.
 - For characterization in terms of spin I: $SO(3) \subset U(d)$

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 $\begin{array}{ll} \mbox{Simple examples} \\ U(2): & \{2,1\} \to SO(3): I = 1/2 \\ U(3): & \{3,1,1\} \to SO(3): I = 2,0 \end{array}$

There exist general rules for $U(d) \rightarrow SO(3)!$

The tool: Young diagrams for irreducible representations of U(d)

- A partition λ sets up a diagram of boxes: $\{2,1\} = -$
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 - Many particle states span the irreps $\lambda = \{\lambda_1, \lambda_2, \dots \lambda_d\}$
 - \blacktriangleright Total spin quantum number calculated by branching rule $U(d){\rightarrow}$ SO(3)

The tool: Young diagrams for irreducible representations of S_n

Partition $\lambda = \{\lambda_1, \lambda_2, \dots, \lambda_p\}$ with $\sum \lambda_i = n$ and all λ_i descending integers describes irrep of \mathbf{S}_n Example \mathbf{S}_3



The answer

Schur-Weyl duality: Irreducible representations of same shapes correlate

Assumption Hilbertspace: $\mathcal{H}_d^N := \mathcal{H}_d \otimes \mathcal{H}_d \otimes \ldots \otimes \mathcal{H}_d$ (*N* copies) with: d = 2I + 1 and *N* number of particles

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The Schur-Weyl theorem

Irreducible reperesentation of the product symmetry group $U(d) \times S_n$:

$$\Gamma_{d}^{(N)} = \bigoplus_{\lambda \vdash N}^{l(\lambda) \leq d} ((\lambda), \{\lambda\})$$
(λ): Irrep of \mathbf{S}_{n} $l(\lambda) \leq d$: Fixed number of rows
(λ): Irrep of $U(d)$ $\lambda \vdash N$: $\sum \lambda_{i} = N$

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"Sum" of identical Young diagrams form irrep for product group

What is it good for?

Simplification of calculations: Example H_5^+

$$\Gamma_2^{(5)} = ((5), \{5\}) \oplus ((41), \{41\}) \oplus ((32), \{32\})$$

The correspondence to usual names							
Young diagram	U(2)	SO(3)	S 5	label			
	$\{5,0\}$	5/2	(5,0)	A_1			
	$\{4,1\}$	3/2	(4,1)	G_1			
	$\{3,2\}$	1/2	(3,2)	H_1			

 $\Gamma_{\rm nspin}(H_5^+) = (6A_1, \mathcal{D}_{5/2}) \oplus (4G_1, 4\mathcal{D}_{3/2}) \oplus (2H_1, 5\mathcal{D}_{1/2})$

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Second part

Statistical state-to-state reaction rates — Application of Schur-Weyls duality

Spin-dependent reaction rates for molecules

The well-known ortho/para species of molecular hydrogen



- No electromagnetic transition between orthoand para states
- Collisional cooling cannot flip spins
- \Rightarrow Reactive collisions?

Spin-flip probability in reactive collisions

Reactive collision Two reactants form intermediate complex before dissociating into initial parts

 Symmetry of the intermediate complex highly influential for reaction rates



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- Reaction rate Probability for one initial state to end up in defined product state
 - Statistics No energy assumptions, only spin-statistical weights and spin conservation
 - A first example: $H_2 + H^+ \rightarrow H_3^+ \rightarrow H_2 + H^+$



The H_5^+ cation

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First results:

Full S_5 MS group: $P(\text{ortho-to-para } H_2) = 9/50$ Smaller group, only H^+ released: $P(\text{ortho-to-para } H_2) = 4/135$

Wrap-up: Schur-Weyl, Young, and reactive collisions

First part: Representations of the nuclear spin wave functions Schur-Weyl duality Prescription for correlation of spin and permutation symmetry

▶ No restriction on I or *N*, easy to implement

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Second part: Reactive collisions

State-to-state rates Specific state transitions E.g. ortho-to-para in $H_2 + H^+$ reaction Intermediate complex Symmetry decisive for reaction paths

Thank you for your attention!



And greetings from Cologne

Schmiedt, Jensen, Schlemmer; JCP **145**, 074301, *(2016)*; Doi: 10.1063/1.4960956